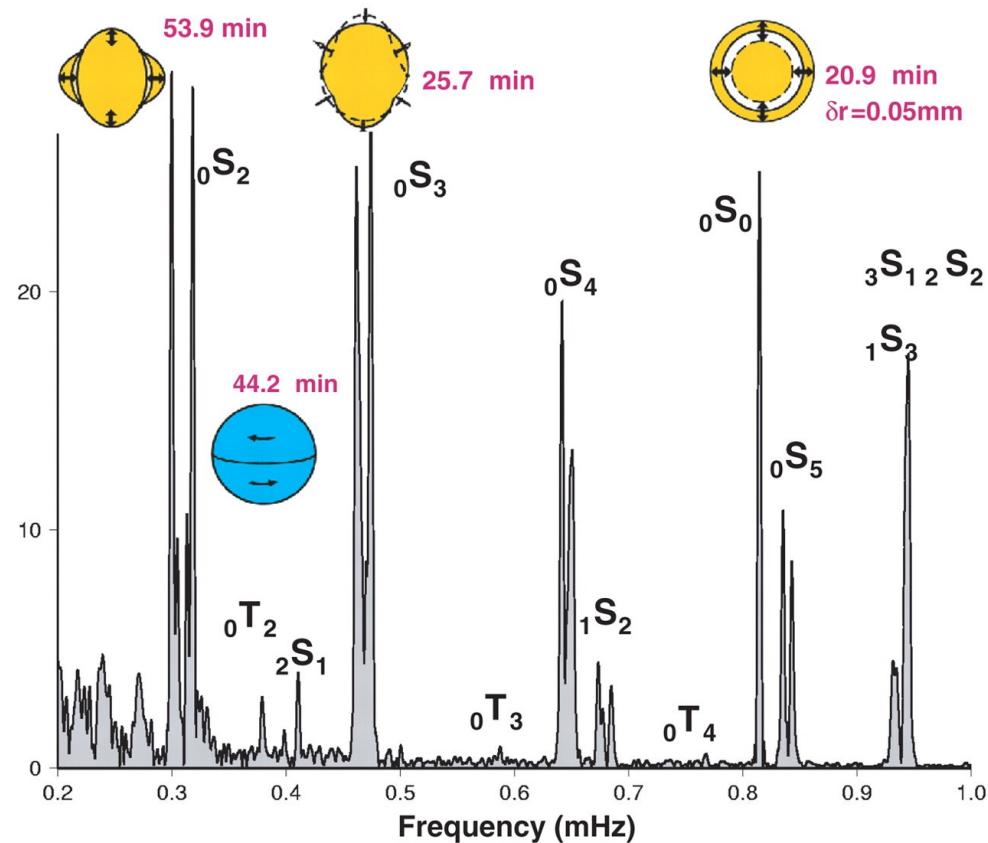


Earth's Free Oscillations (Normal Modes)

TA Lecture for GP 238, Spring 2023

Ji, Qing

May 18



Cover figure from:

Park *et al.* (2005). Earth's Free Oscillations Excited by the 26 December 2004 Sumatra-Andaman Earthquake. *Science*.

Our Earth can vibrate as a whole, just like a ringing bell!

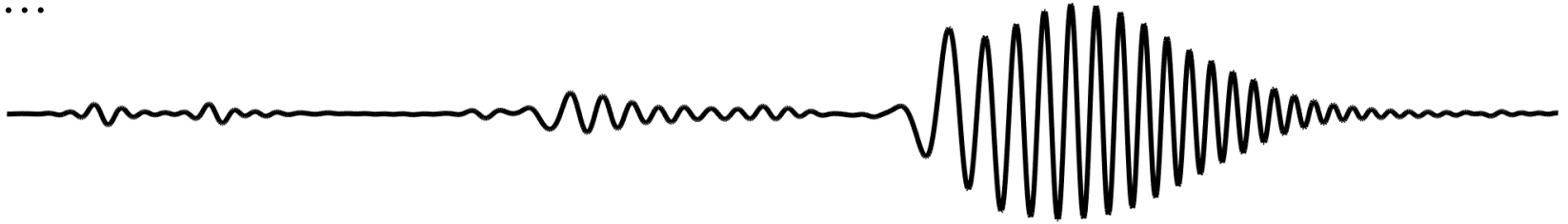
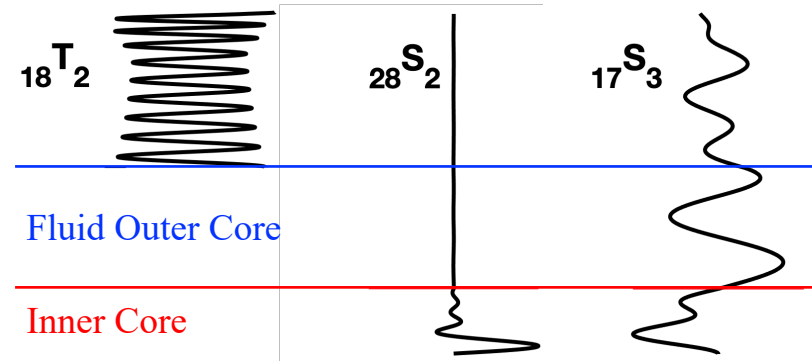


After large earthquakes or explosions, free oscillations are observed just like the bell after being struck

With much exaggeration of course

These normal modes can be applied to

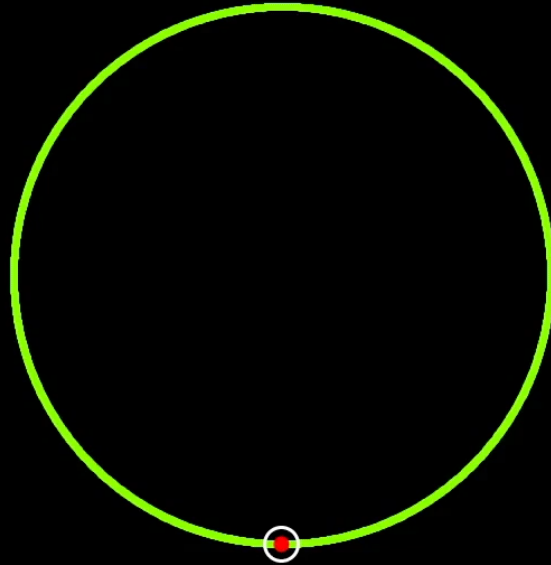
- study the (deep) interior of the Earth
- compute synthetic seismic waveforms
- provide info on earthquake size & duration
-



Traveling wave ----> Standing wave

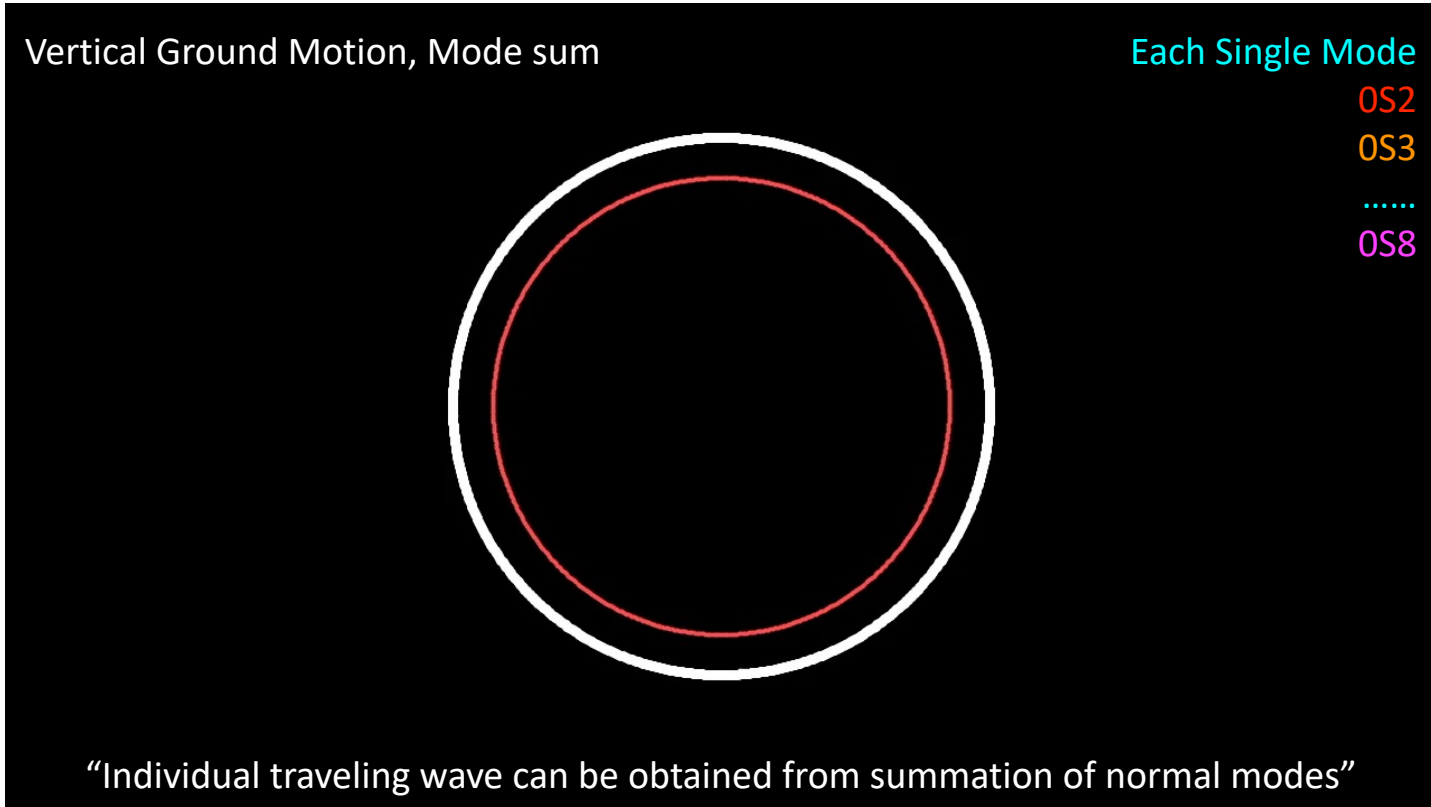
Vertical & Radial Ground Motion

Synthetic seismic wavefield



“Individual normal mode can be obtained from superposition of travelling waves”

Standing wave ----> Traveling wave

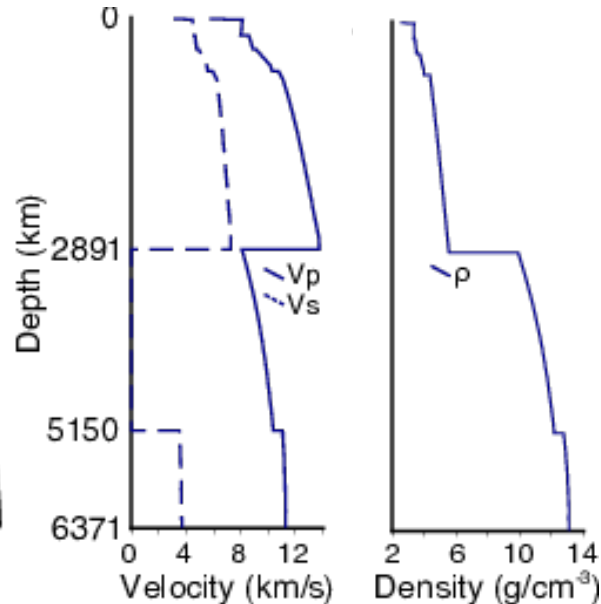
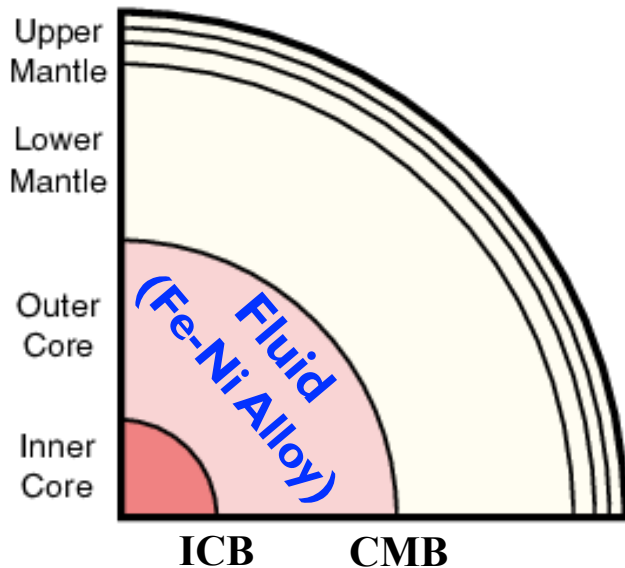


Normal Modes in SNREI Earth Model

SNREI: **S**pherically symmetric, **N**on-**R**otating, **E**lastic and **I**sotropic

Material properties only vary in radial direction

Preliminary Reference Earth Model (PREM)



Dziewonski & Anderson (1981)

Obtained by joint inversion of:

- ✓ Body wave travel times
- ✓ Periods of **normal modes** (~ 900 modes)

Figure from CIDER

In this lecture

Traveling wave & Standing wave: **Ray-mode** duality

- ✓ Individual **normal mode** can be obtained from **superposition** of **travelling waves**
- ✓ Individual **traveling wave** (P, S, surface waves ...) can be obtained from **summation** of **normal modes**

We can get a sense of this duality by visualizing radial eigenfunction

Excitation of normal modes: Calculate Green's function, synthetic seismogram

(Complete Basis!)

$$\mathbf{G}(\mathbf{r}, t; \mathbf{r}_0) = \sum_{n=0}^{\infty} \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{\sin(n\omega_l t)}{n\omega_l} {}_n\mathbf{S}_{lm}(\mathbf{r}) {}_n\mathbf{S}_{lm}(\mathbf{r}_0), \quad t > 0.$$

Procedure to solve normal modes

- Linearized equations of motion in frequency domain, SNREI model

$$-\omega^2 \rho \mathbf{u} - \nabla \cdot \boldsymbol{\sigma} + \rho \nabla(gu_r) + \rho \nabla \phi - \rho (\nabla \cdot \mathbf{u}) g \hat{\mathbf{r}} = 0$$

Hooke's law:

$$\boldsymbol{\sigma} = K(\nabla \cdot \mathbf{u})\mathbf{I} + 2\mu \left[\frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^T) - \frac{1}{3}(\nabla \cdot \mathbf{u})\mathbf{I} \right]$$

Gravitational effect is important at low frequency

- Linearized boundary conditions

	Solid-Solid	Solid-Fluid	Free surface
Kinematic (Displacement \mathbf{u})	$[\mathbf{u}]_{\pm}^{\pm} = \mathbf{0}$	$[\hat{\mathbf{n}} \cdot \mathbf{u}]_{\pm}^{\pm} = 0$	————
Dynamic (Normal traction $\mathbf{T} = \hat{\mathbf{n}} \cdot \boldsymbol{\sigma}$)	$[\mathbf{T}]_{\pm}^{\pm} = \mathbf{0}$	$\hat{\mathbf{n}}[\hat{\mathbf{n}} \cdot \mathbf{T}]_{\pm}^{\pm} = [\mathbf{T}]_{\pm}^{\pm} = \mathbf{0}$	$\mathbf{T} = \mathbf{0}$

Procedure to solve normal modes

- Expand displacement field \mathbf{u} with Vector Spherical Harmonics. For each mode:

$${}_n\mathbf{u}_{lm}(r, \theta, \phi) = {}_nU_l(r)\mathbf{P}_{lm}(\theta, \phi) + {}_nV_l(r)\mathbf{B}_{lm}(\theta, \phi) + {}_nW_l(r)\mathbf{C}_{lm}(\theta, \phi)$$

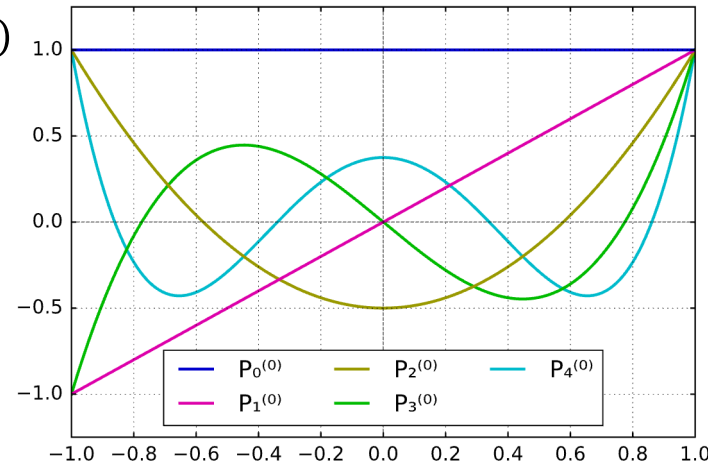
SNREI model leads to **degeneracy in m** . Mode frequency and radial function only depend on n and l

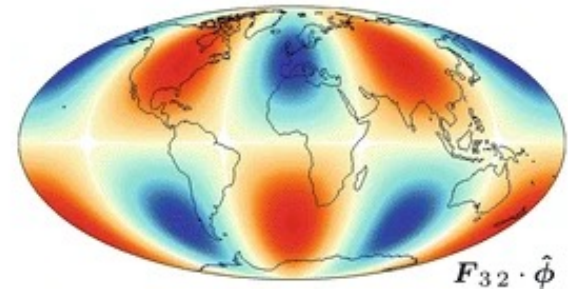
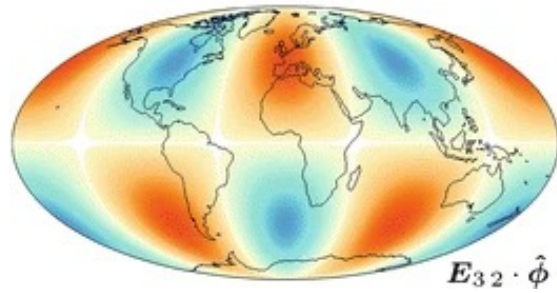
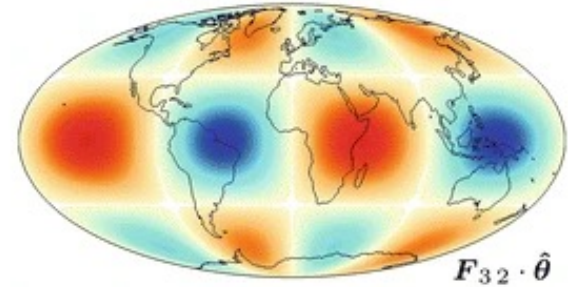
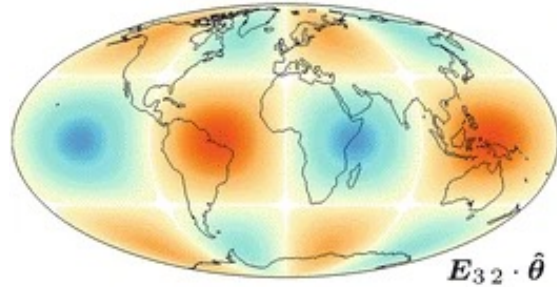
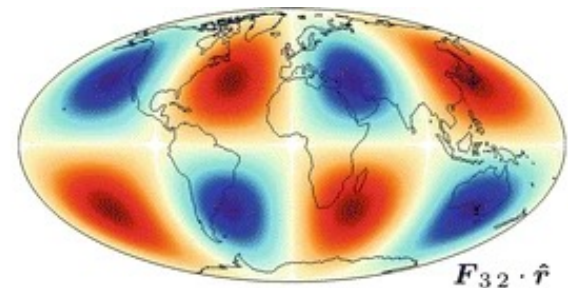
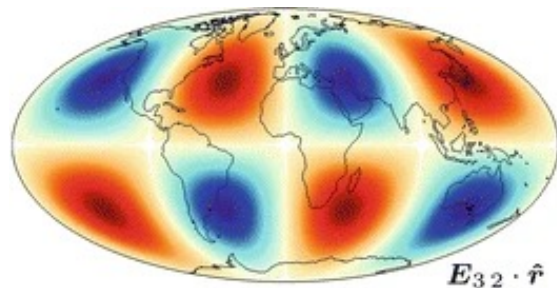
Vector Spherical Harmonics: $\mathbf{P}_{lm}(\theta, \phi)$, $\mathbf{B}_{lm}(\theta, \phi)$, $\mathbf{C}_{lm}(\theta, \phi)$

- Related to real-valued spherical harmonics $\psi_{lm}(\theta, \phi)$ and its surface gradient and surface curl
- Product between $P_{lm}(\cos \theta)$ and $\sin(m\phi)$ or $\cos(m\phi)$

n, l, m indicate number of nodes in the eigenfunctions for radial, latitudinal and longitudinal directions

Plots of associated Legendre polynomials from Wikipedia





n, l, m are related to number of nodes
in the eigenfunctions for radial,
latitudinal and longitudinal directions

Higher orders indicate smaller
wavelength, finer spatial resolution

Procedure to solve normal modes

- Substitute the expansion into wave equation
- Decouple into **P-SV (spheroidal)** and **SH (toroidal)** systems
- Change problem from a second-order ODE into first-order ODE system by introducing traction vector

$${}_n \mathbf{u}_{lm}(r, \theta, \phi) = {}_n U_l(r) \mathbf{P}_{lm}(\theta, \phi) + {}_n V_l(r) \mathbf{B}_{lm}(\theta, \phi) + {}_n W_l(r) \mathbf{C}_{lm}(\theta, \phi)$$

$$\hat{\mathbf{r}} \cdot \boldsymbol{\sigma} = {}_n \mathbf{T}_{lm}(r, \theta, \phi) = \underbrace{{}_n R_l(r) \mathbf{P}_{lm}(\theta, \phi) + {}_n S_l(r) \mathbf{B}_{lm}(\theta, \phi)}_{\text{P-SV}} + \underbrace{{}_n T_l(r) \mathbf{C}_{lm}(\theta, \phi)}_{\text{SH}}$$

P-SV

SH

Example: ODEs for SH (toroidal) system

$$r^{-2} \frac{d}{dr} \left[\mu r^2 \left(\frac{dW}{dr} - r^{-1} W \right) \right] + \mu r^{-1} \left(\frac{dW}{dr} - r^{-1} W \right) + [\omega^2 \rho - (l^2 + l - 2) \mu r^{-2}] W = 0$$

$$T = \mu \left(\frac{dW}{dr} - r^{-1} W \right)$$

$$\frac{dW}{dr} = r^{-1} W + \mu^{-1} T$$

$$\frac{dT}{dr} = [-\omega^2 \rho + (l^2 + l - 2) \mu r^{-2}] W - 3r^{-1} T$$

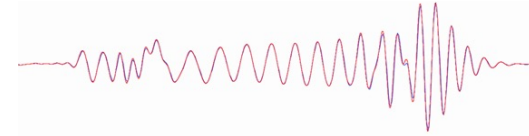
Procedure to solve normal modes

$$\frac{dW}{dr} = r^{-1}W + \mu^{-1}T$$

$$\frac{dT}{dr} = [-\omega^2\rho + (l^2 + l - 2)\mu r^{-2}]W - 3r^{-1}T$$

Mineos

User Manual
Version 1.0.2



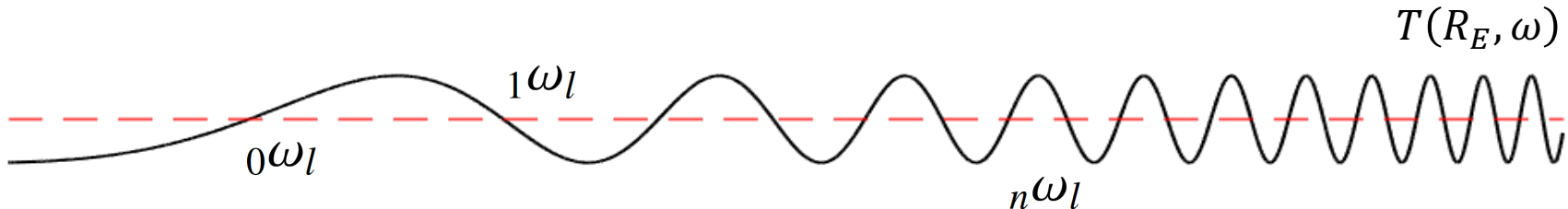
➤ Solve this eigenvalue problem (for toroidal modes)

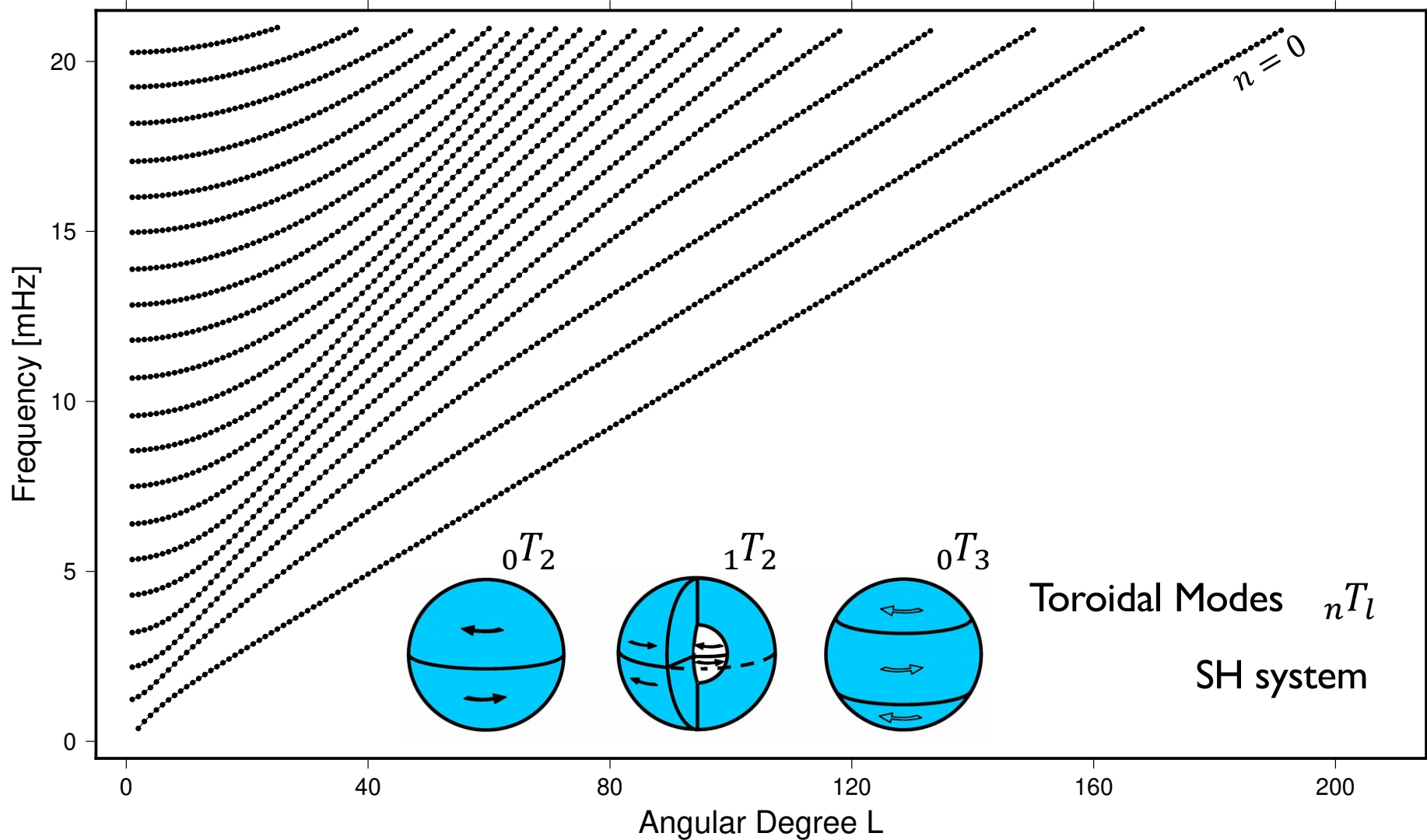
- Choose l and a trial frequency ω
- Start from the Core-Mantle Boundary (CMB, fluid-solid interface, 2891 km depth)

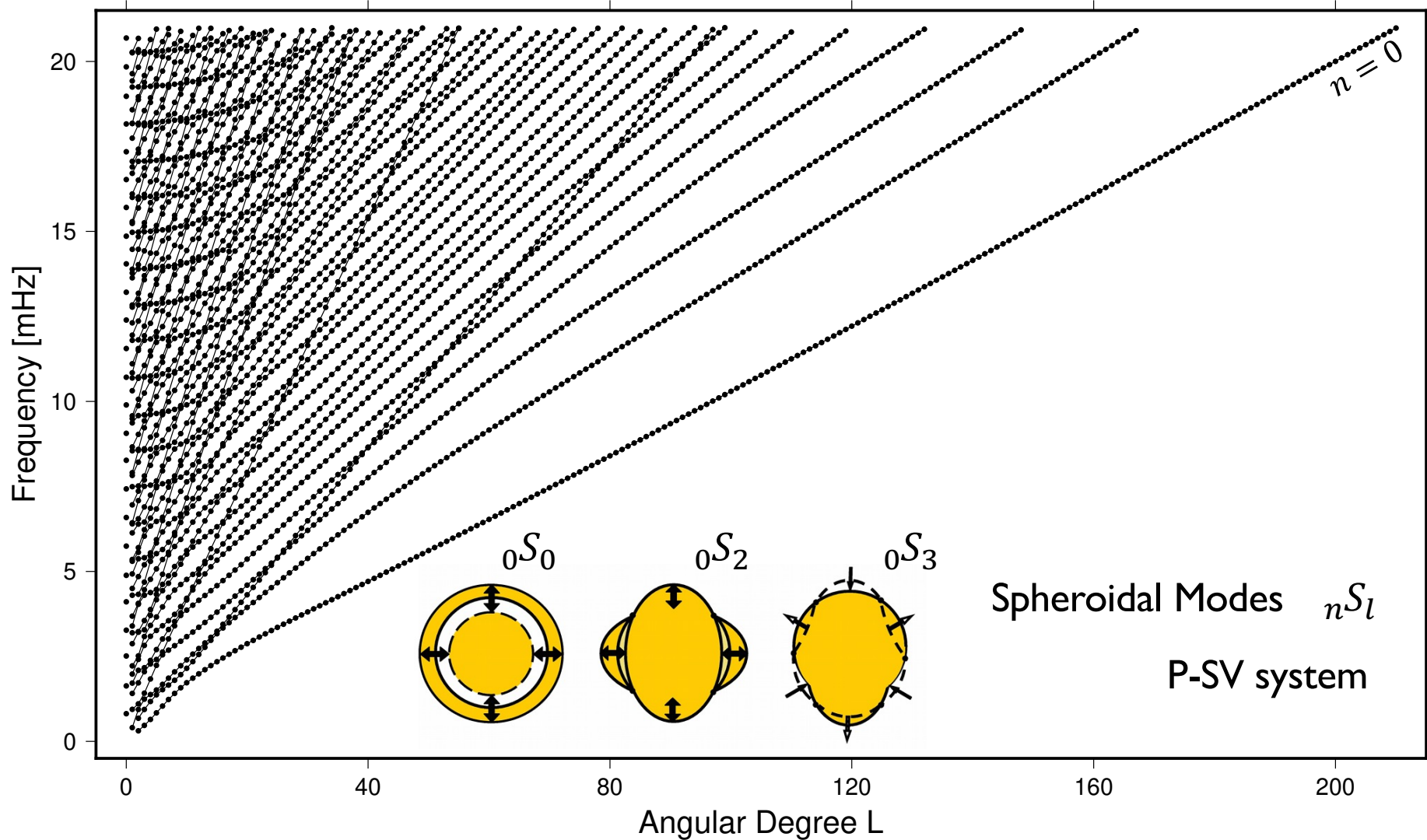
$$W(r = \text{CMB}) = \text{any nonzero} \rightarrow 1$$

$$T(r = \text{CMB}) = 0$$

- Integrate upward until reaching the free surface
- Satisfying the free-surface boundary condition gives constraint on mode frequency





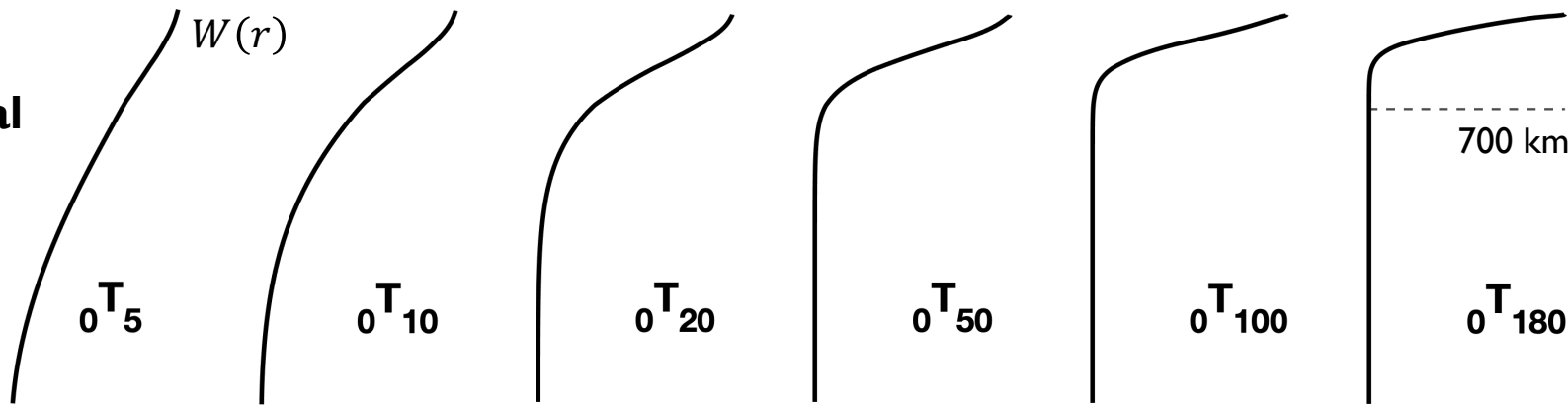


Radial Eigenfunction of Toroidal Modes

$${}_n W_l(r) C_{lm}(\theta, \phi)$$

$n = 0$

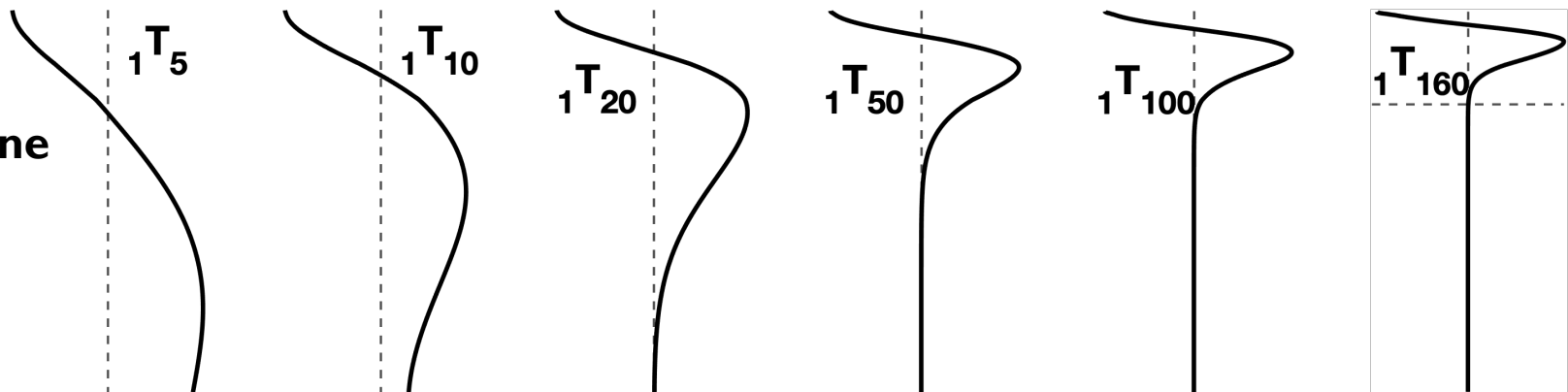
**Fundamental
Love Mode**



$n = 1$

1st Overtone

1 node

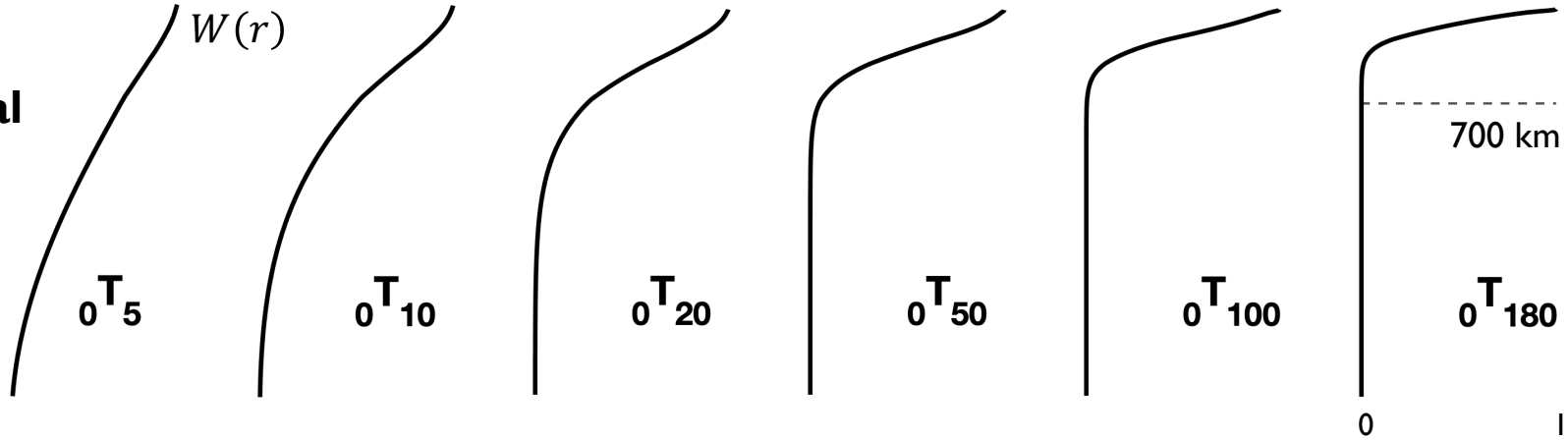


Excitation of surface waves

$${}_n W_l(r) \mathbf{C}_{lm}(\theta, \phi)$$

$n = 0$

**Fundamental
Love Mode**



Excitation “coefficient”: Radial eigenfunction evaluated at source depth

$$\mathbf{G}(\mathbf{r}, t; \mathbf{r}_0) = \sum_{n=0}^{\infty} \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{\sin(n\omega_l t)}{n\omega_l} {}_n \mathbf{S}_{lm}(\mathbf{r}) \boxed{{}_n \mathbf{S}_{lm}(\mathbf{r}_0)}, \quad t > 0.$$

Source with depth h generates surface waves with wavelength $\lambda \gtrsim h$

Simply speaking, shallow earthquakes can efficiently generate surface waves

Group velocity dispersion of Love wave (at long periods)

Group Velocity [km/s]

8

6

4

0

5

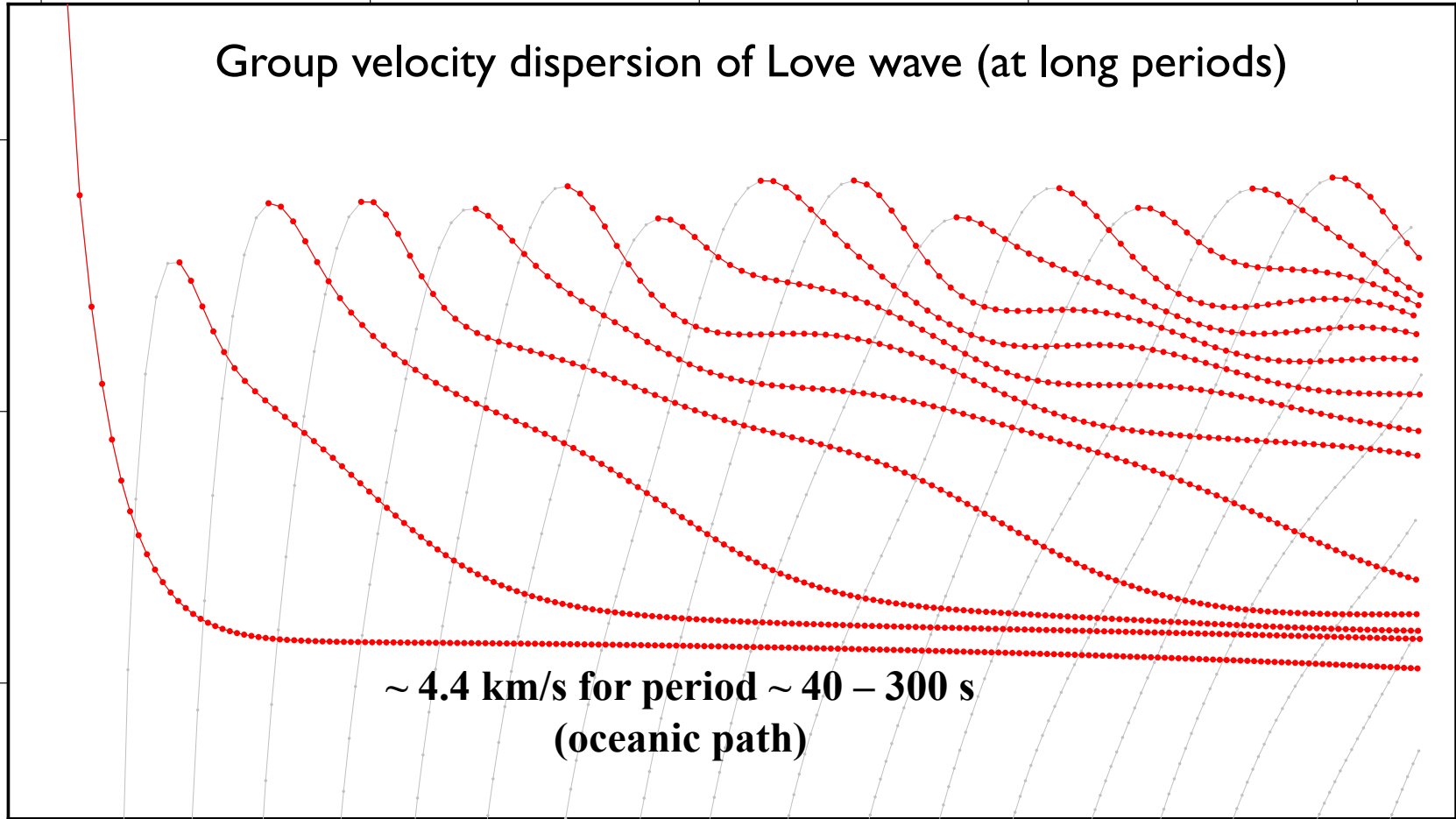
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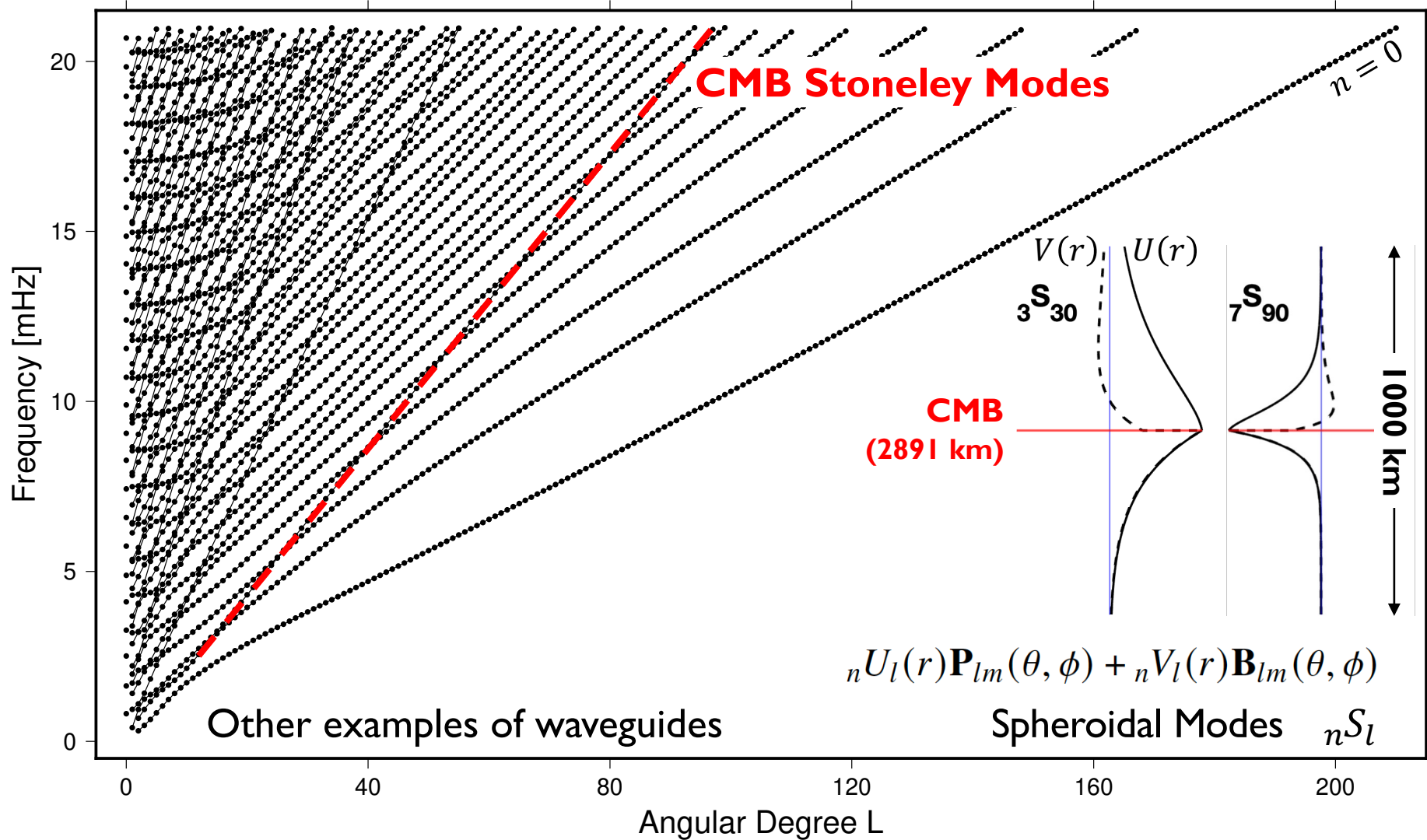
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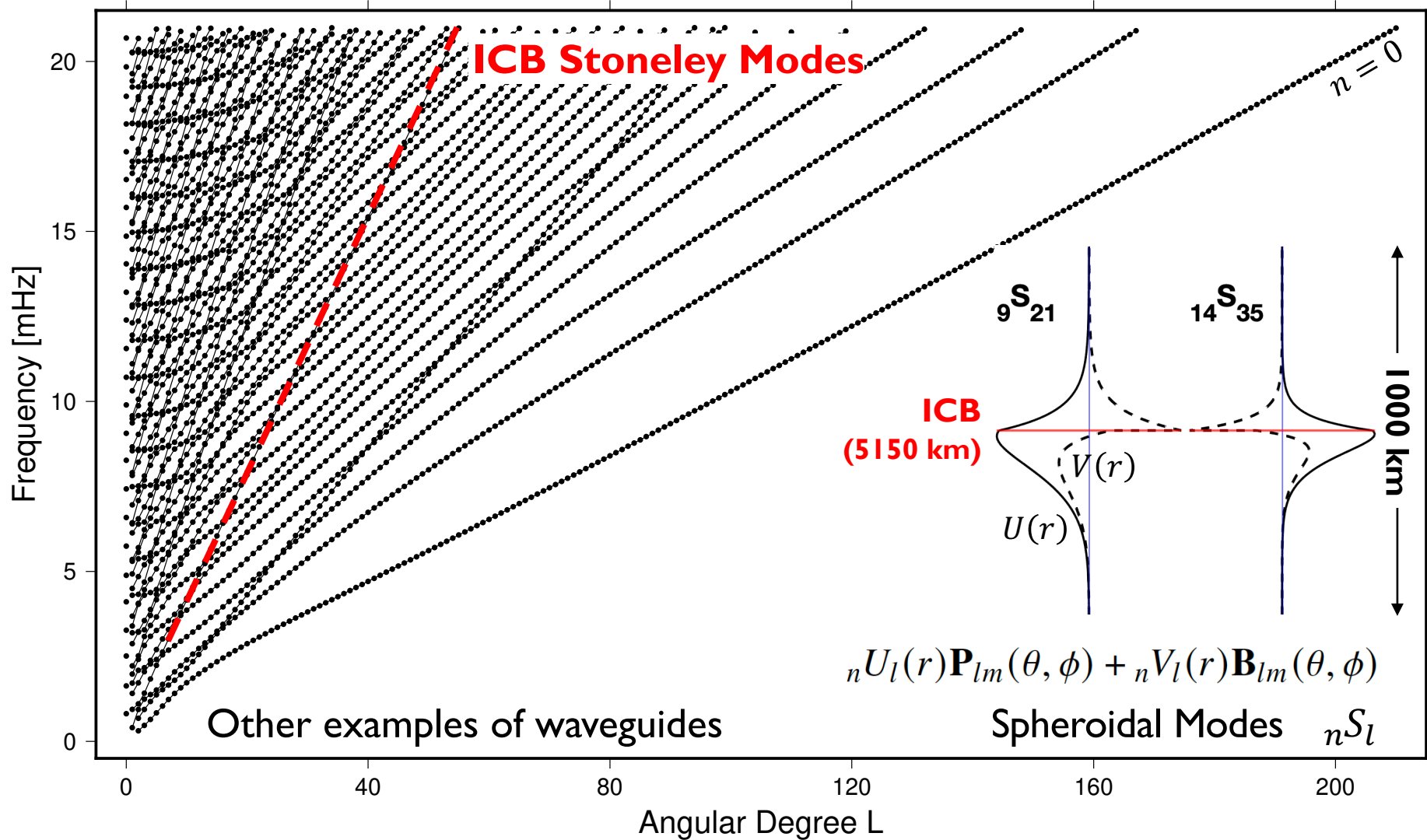
20

Frequency [mHz]

**~ 4.4 km/s for period $\sim 40 - 300$ s
(oceanic path)**



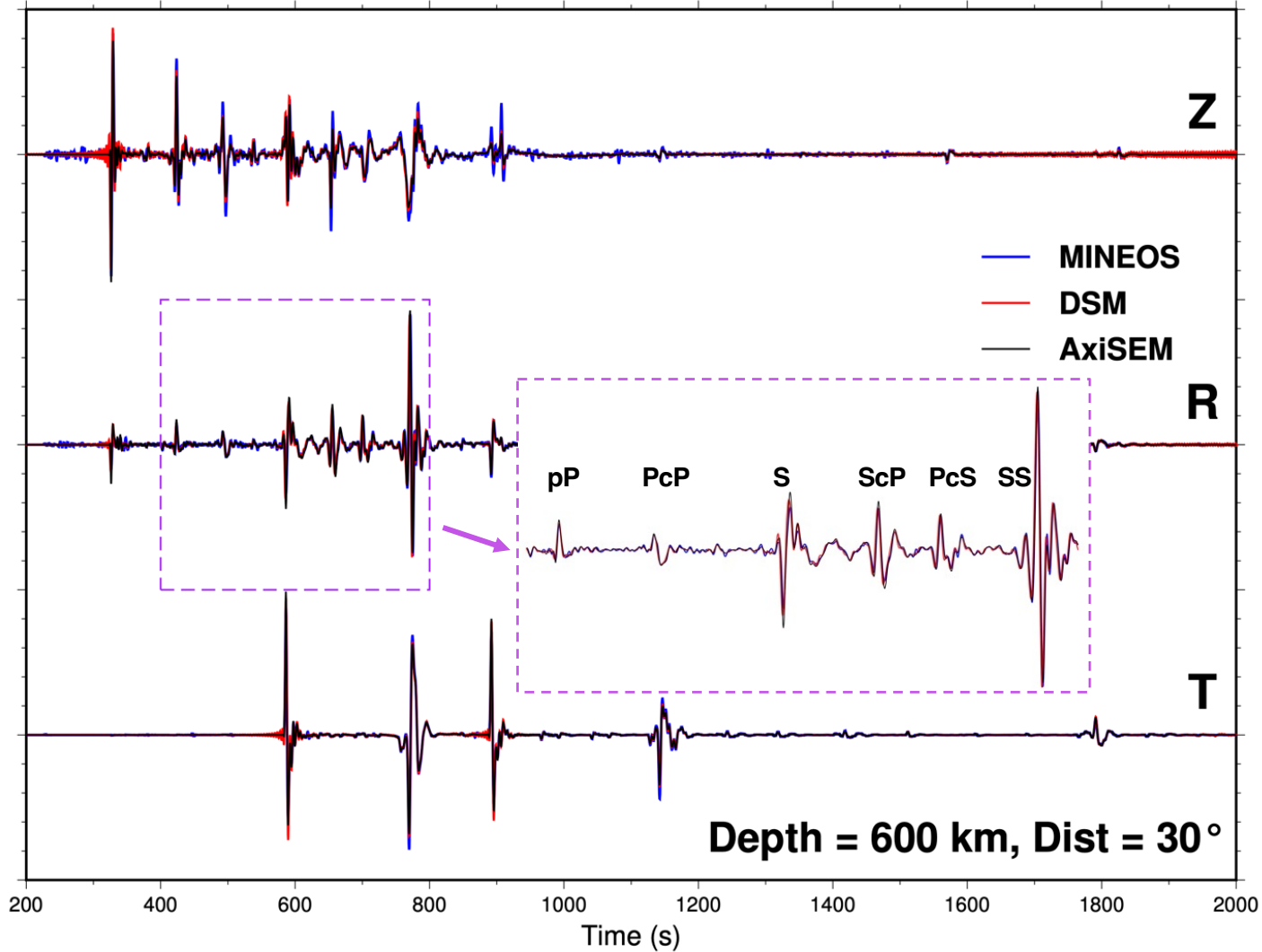




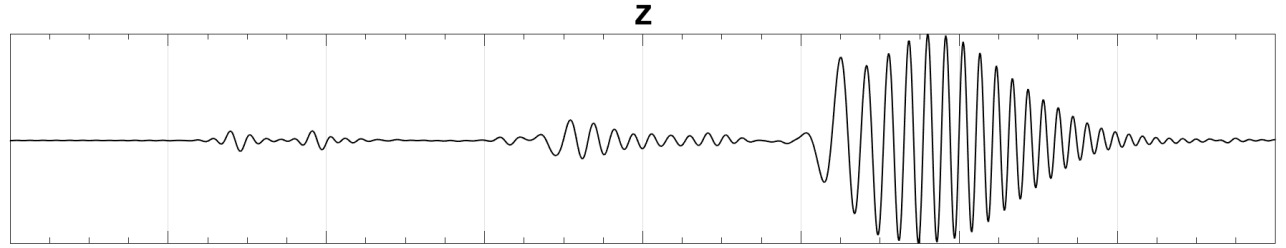
Calculate synthetic
seismograms from
mode summation

Deep earthquake
(Depth 600 km)

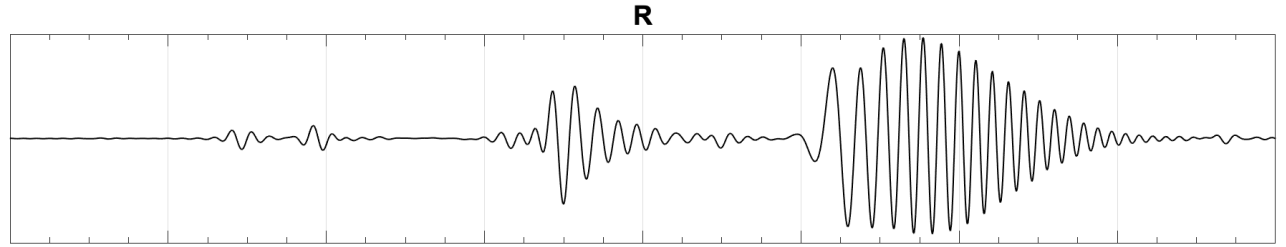
Low-pass 0.2 Hz



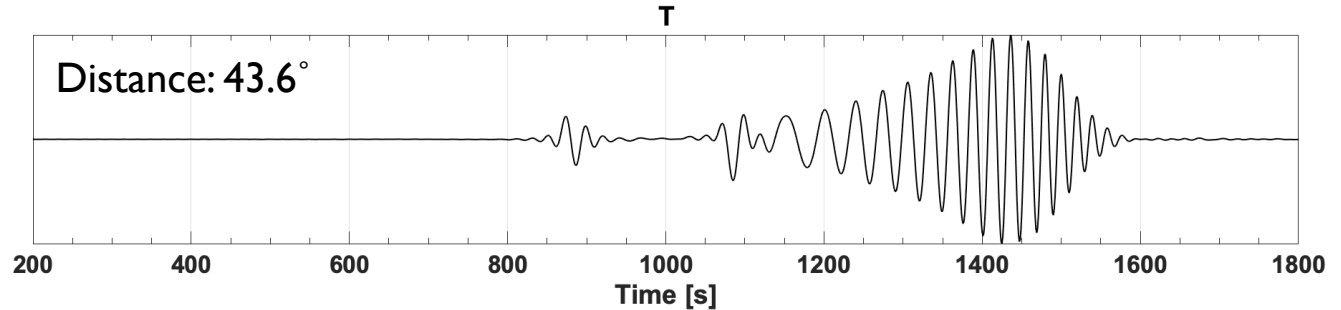
Calculate synthetic
seismograms from
mode summation



Shallow earthquake
(Depth 10 km)



Low-pass 0.05 Hz



In this lecture

Traveling wave & Standing wave: **Ray-mode** duality

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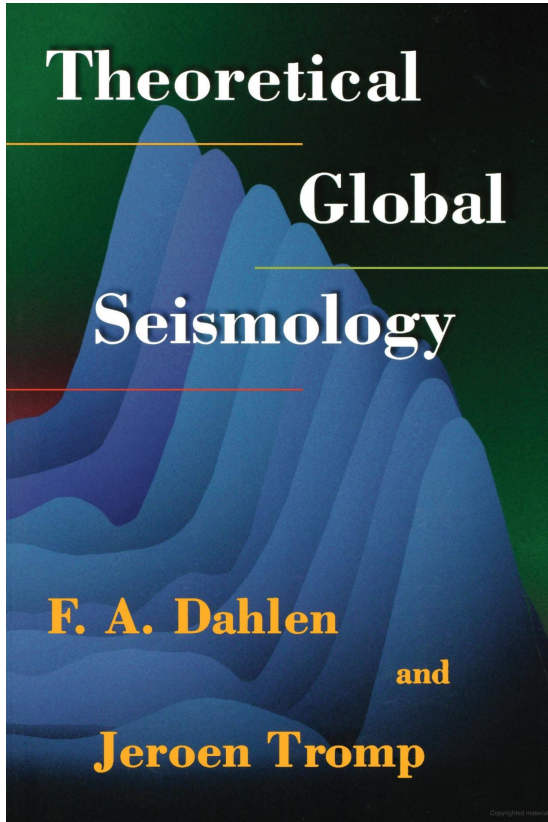
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Recommended Reference (if you are interested!)



Chapter 8.7 & 8.8: Gallery of radial eigenfunctions

Chapter 10: Synthetic seismograms from mode summation

Singh, S.J., Rani, S. (2020). Free Oscillations of the Earth.
In: Encyclopedia of Solid Earth Geophysics.

Thank You!